

Andreas Ioannidis, Linnæus University, Sweden

The Eigenvalue Problem for the Cavity Maxwell Operator

Abstract. In this talk we discuss the propagation problem of the time harmonic electromagnetic field $\mathbf{e} := (\mathbf{E}, \mathbf{H})^T$ (time convention is taken to be $e^{-i\omega t}$) inside a Lipschitz bounded domain $\Omega \subset \mathbb{R}^3$. Ω represents a cavity, which is filled with a general bianisotropic medium and is perfect conducting, that is, the boundary condition $\hat{\mathbf{n}} \times \mathbf{E} = \mathbf{0}$ on $\partial\Omega$ is satisfied. $\hat{\mathbf{n}}$ denotes the outward normal, which is defined almost everywhere on $\partial\Omega$.

This problem is formally formulated as an eigenvalue problem for a linear operator pencil

$$\mathcal{Q}\mathbf{e} = \omega\mathbf{M}(\omega)\mathbf{e}. \quad (\star)$$

Here \mathcal{Q} stands for the self-adjoint Maxwell operator

$$\mathcal{Q} := i \begin{bmatrix} 0 & -\text{curl} \\ \text{curl} & 0 \end{bmatrix},$$

and $\mathbf{M} = \mathbf{M}(\omega)$ for the material matrix

$$\mathbf{M} := \begin{bmatrix} \boldsymbol{\varepsilon} & \boldsymbol{\xi} \\ \boldsymbol{\zeta} & \boldsymbol{\mu} \end{bmatrix},$$

which models the medium properties inside Ω . The entries of \mathbf{M} are $L^\infty(\Omega)$ functions that depend on the eigenvalue ω .

We realize (\star) in a Hilbert space setting and our analysis breaks down into two steps:

1. We first study the hollow cavity, modeled by the simplest case $\mathbf{M} = I$. Then (\star) becomes a standard eigenvalue problem for which we employ a variant of the spectral theorem for discrete self-adjoint operators to prove existence for the eigenelements.
2. For the general matrix operator \mathbf{M} , we use perturbation arguments to prove existence for the eigenelements of the general cavity and characterize them with aim of their hollow cavity counterparts.