

Spectral Theory for Operator Matrices Related to Models in Mechanics

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Abstract

We derive various properties of the operator matrix $\mathcal{A} = \begin{bmatrix} 0 & I \\ -A_0 & -D \end{bmatrix}$, where A_0 is a uniformly positive operator and $A_0^{-1/2}DA_0^{-1/2}$ is a bounded non-negative operator in a Hilbert space H . Such operator matrices are associated with second order problems of the form $\ddot{z}(t) + A_0z(t) + D\dot{z}(t) = 0$ which are used as models for transverse motions of thin beams in the presence of damping.

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Key words: Operator matrices, second order equations, spectrum, Riesz basis, definitizable operator, Krein space, analytic semigroup.

1 Introduction

A linear equation describing transverse motions of a thin beam can be written in the form

$$\frac{\partial^2 u}{\partial t^2} + \frac{\partial^2}{\partial r^2} \left[E \frac{\partial^2 u}{\partial r^2} + C_d \frac{\partial^3 u}{\partial r^2 \partial t} \right] = 0, \quad r \in (0, 1), t > 0,$$

where $u(r, t)$ is the transverse displacement of the beam at time t and position r . The existence and behaviour of solutions u depend also on boundary and initial conditions. In the example above we are interested in solutions having finite energy, i.e. solutions such that $\|u(\cdot, t)\|^2 + \|u''(\cdot, t)\|^2 < \infty$ for all $t > 0$ where $\|\cdot\|$ denotes the usual norm in the Hilbert space $L^2(0, 1)$. Identifying the function $u(\cdot, t)$ with an element $z(t) \in L^2(0, 1)$ by $z(t)(r) = u(r, t)$ we obtain from the partial differential equation above a second order equation in $L^2(0, 1)$ of the form

$$\ddot{z}(t) + A_0z(t) + D\dot{z}(t) = 0, \tag{1}$$

where $A_0 = E \frac{\partial^4}{\partial r^4}$, $D = \frac{\partial^2}{\partial r^2} C_d \frac{\partial^2}{\partial r^2}$ acting in $L^2(0, 1)$ with appropriate domains encoding the boundary conditions under consideration.

In this paper we study second order equations of type (1) in an abstract Hilbert space H where the stiffness operator A_0 is a possibly unbounded positive operator on H and is assumed to be boundedly invertible, and D , the damping operator, is an unbounded operator on H . This second order equation is equivalent to the standard first-order equation $\dot{x}(t) = \mathcal{A}x(t)$ where $\mathcal{A} : \mathcal{D}(\mathcal{A}) \subset \mathcal{D}(A_0^{1/2}) \times H \rightarrow \mathcal{D}(A_0^{1/2}) \times H$ is given by

$$\mathcal{A} = \begin{bmatrix} 0 & I \\ -A_0 & -D \end{bmatrix},$$

$$\mathcal{D}(\mathcal{A}) = \left\{ \begin{pmatrix} z \\ w \end{pmatrix} \in \mathcal{D}(A_0^{1/2}) \times \mathcal{D}(A_0^{1/2}) \mid A_0 z + Dw \in H \right\}.$$

This operator matrix has been studied in the literature for more than 20 years. Interest in this particular model is motivated by various problems such as stabilization, see for example [7], [29], [30], [32], solvability of Riccati equations [14], minimum-phase property [20] and compensator problems with partial observations [15]. It is well-known that \mathcal{A} generates a C_0 -semigroup of contractions in $\mathcal{D}(A_0^{1/2}) \times H$, where $\mathcal{D}(A_0^{1/2})$ is equipped with the norm $x \mapsto \|A_0^{1/2}x\|_H$, and thus the spectrum of \mathcal{A} is located in the closed left half plane. This goes back to [4] and [28], see also [5], [9]. Several authors have proved independently of each other that the condition

$$\inf_{z \in \mathcal{D}(A_0^{1/2}) \setminus \{0\}} \frac{\langle A_0^{-1/2} D z, A_0^{1/2} z \rangle_H}{\|z\|_H^2} > 0$$

is sufficient for exponential stability of the C_0 -semigroup generated by \mathcal{A} , see for example [4], [5], [6], [9], [16], [17], [34] and [35]. Other properties of the C_0 -semigroup such as analyticity have been studied in [4], [5], [9], [10], [11], [12], [13], [18], [19] and [25].

Most of the papers require that the damping operator D is comparable with A^ρ for some $\rho \in [1/2, 1]$. In [25] the damping operator D is of the form

$$D = \alpha A_0 + B, \tag{2}$$

where $\alpha > 0$ is a constant, A_0^{-1} is compact and B is symmetric and A_0 -compact. If $-1/\alpha \notin \sigma_p(\mathcal{A})$, then it is shown in [25] that \mathcal{A} generates an analytic semigroup. In [22] we extend this result to the case that A_0^{-1} is compact in H and $0 \notin \sigma_{ess}(A_0^{-1}D)$. Moreover, it was shown in [22] that \mathcal{A} is a definitizable operator.

In this paper we also assume that A_0^{-1} is a compact operator. We show that the condition $\sigma_{ess}(A_0^{-1}D) = \{0\}$ is equivalent to the definitizability of \mathcal{A} . It turns out, that in the case that \mathcal{A} has a compact resolvent, in contrast to the situation considered in [25] and [22], the operator \mathcal{A} is no longer definitizable.

Throughout this paper we assume that all Hilbert spaces are infinite dimensional.

2 Spectrum of operators in Krein spaces

Let $(\mathcal{H}, [\cdot, \cdot])$ be a Krein space. We briefly recall that a complex linear space \mathcal{H} with a hermitian nondegenerate sesquilinear form $[\cdot, \cdot]$ is called a *Krein space* if there exists a decomposition $\mathcal{H} = \mathcal{H}_+ \oplus \mathcal{H}_-$ with subspaces \mathcal{H}_\pm being orthogonal to each other with respect to $[\cdot, \cdot]$ such that $(\mathcal{H}_\pm, \pm[\cdot, \cdot])$ are Hilbert spaces. In the following all topological notions are understood with respect to some Hilbert space norm $\|\cdot\|$ on \mathcal{H} such that $[\cdot, \cdot]$ is $\|\cdot\|$ -continuous. For the basic theory of Krein spaces and operators acting therein we refer to [8] and [2].

Let A be a closed operator in \mathcal{H} . We define the extended spectrum $\sigma_e(A)$ of A by $\sigma_e(A) := \sigma(A)$ if A is bounded and $\sigma_e(A) := \sigma(A) \cup \{\infty\}$ if A is unbounded. The resolvent set of A is denoted by $\rho(A)$ and we set $\rho_e(A) := (\mathbb{C} \cup \{\infty\}) \setminus \sigma_e(A)$. A point $\lambda_0 \in \mathbb{C}$ is said to belong to the *approximative point spectrum* $\sigma_{ap}(A)$ of A if there exists a sequence $(x_n) \subset \mathcal{D}(A)$ with $\|x_n\| = 1$, $n = 1, 2, \dots$, and $\|(A - \lambda_0 I)x_n\| \rightarrow 0$ if $n \rightarrow \infty$. For a self-adjoint operator A in \mathcal{H} all real spectral points of A belong to $\sigma_{ap}(A)$ (see e.g. [8, Corollary VI.6.2]).

The indefiniteness of the scalar product on \mathcal{H} leads to the definition of several subsets of the spectrum of an operator. The following definition was given in [24], [27] and [3].

Definition 2.1 *For a self-adjoint operator A in \mathcal{H} with domain $\mathcal{D}(A)$ a point $\lambda_0 \in \sigma(A)$ is called a spectral point of type π_+ (type π_-) of A if $\lambda_0 \in \sigma_{ap}(A)$ and if there exists a linear submanifold $\mathcal{H}_0 \subset \mathcal{H}$ with $\text{codim } \mathcal{H}_0 < \infty$ such that for every sequence $(x_n) \subset \mathcal{H}_0 \cap \mathcal{D}(A)$ with $\|x_n\| = 1$ and $\|(A - \lambda_0 I)x_n\| \rightarrow 0$ as $n \rightarrow \infty$ we have*

$$\liminf_{n \rightarrow \infty} [x_n, x_n] > 0 \quad (\text{resp. } \limsup_{n \rightarrow \infty} [x_n, x_n] < 0).$$

A point $\lambda_0 \in \sigma(A)$ is called a spectral point of positive type (negative) type of A if $\lambda_0 \in \sigma_{\pi_+}(A)$ (resp. $\lambda \in \sigma_{\pi_-}(A)$) and $\mathcal{H}_0 = \mathcal{H}$. The point ∞ is said to

be of positive (negative) type of A if A is unbounded and for every sequence $(x_n) \subset \mathcal{D}(A)$ with $\lim_{n \rightarrow \infty} \|x_n\| = 0$ and $\|Ax_n\| = 1$ we have

$$\liminf_{n \rightarrow \infty} [Ax_n, Ax_n] > 0 \quad (\text{resp. } \limsup_{n \rightarrow \infty} [Ax_n, Ax_n] < 0).$$

We denote the set of all points of $\sigma_e(A)$ of positive (negative) type by $\sigma_{++}(A)$ (resp. $\sigma_{--}(A)$).

Similarly, one introduces the notion that ∞ is a spectral point of type π_+ (type π_-) of A (cf. [3]). We denote the set of all points of $\sigma_e(A)$ of type π_+ (type π_-) of A by $\sigma_{\pi_+}(A)$ (resp. $\sigma_{\pi_-}(A)$). It is not difficult to see that the sets $\sigma_{++}(A)$ and $\sigma_{--}(A)$ are contained in $\overline{\mathbb{R}}$. Moreover the non-real spectrum of A cannot accumulate to $\sigma_{++}(A) \cup \sigma_{--}(A)$. More properties of these sets can be found in [3].

Recall that a self-adjoint operator A in a Krein space $(\mathcal{H}, [\cdot, \cdot])$ is called *definitizable* if $\rho(A) \neq \emptyset$ and if there exists a rational function $p \neq 0$ having poles only in $\rho(A)$ such that $[p(A)x, x] \geq 0$ for all $x \in \mathcal{H}$. Then the spectrum of A is real or its non-real part consists of a finite number of points, cf. [26].

The theorems below will be used in the proof of our main result in Section 3 below. It shows that a real spectral point which is not an eigenvalue and which is adjacent to the spectrum of type π_+ /type π_- and to the resolvent set of a definitizable operator has to be a spectral point of positive/negative type.

Theorem 2.2 *Let A be a definitizable operator and let (a, b) be an open interval such that*

$$(a, b) \subset \sigma_{\pi_+}(A) \cup \sigma_{\pi_-}(A) \cup \rho(A) \quad \text{and} \quad b \notin \sigma_p(A).$$

If there exists $c > b$ with $(b, c) \subset \rho(A)$, then

$$b \in \sigma_{++}(A) \cup \sigma_{--}(A) \cup \rho(A).$$

Proof:

As A is a definitizable operator there exists $a_0 \in (a, b)$ and $c_0 \in (b, c)$ such that $[a_0, c_0] \setminus \{b\}$ is either a subset of $\sigma_{++}(A) \cup \rho(A)$ or a subset of $\sigma_{--}(A) \cup \rho(A)$. Let $[a_0, c_0] \setminus \{b\} \subset \sigma_{++}(A) \cup \rho(A)$. We denote by E the spectral function of A (cf. [26]). Then $E([a_0, c_0])$ is defined and $E([a_0, c_0])\mathcal{H}$ is a Krein space. By [23, Theorem 4.8 and Remark 4.9] the operator

$$A' := A|E([a_0, c_0])\mathcal{H}.$$

is definitizable and the restriction of E to $E([a_0, c_0])$ is the spectral function of A' . By Proposition [3, Proposition 25], for every interval $\delta \subset [a_0, c_0]$ with

$b \notin [a_0, c_0]$ the projection $E([a_0, c_0])$ is non-negative. Therefore, the subspace \mathcal{S} of $E([a_0, c_0])\mathcal{H}$,

$$\mathcal{S} := \text{cls} \{E(\delta)\mathcal{H} : \delta \subset [a_0, c_0], \delta \text{ interval}, b \notin \bar{\delta}\}.$$

is non-negative ([26, Lemma I.5.3]). By ([26, Propositions II.5.1 and II.5.2]) the orthogonal subspace to \mathcal{S} in $E([a_0, c_0])\mathcal{H}$ is the root subspace of A corresponding to b , hence, as $b \notin \sigma_p(A)$,

$$\mathcal{S} = E([a_0, c_0])\mathcal{H},$$

that is, $(E([a_0, c_0])\mathcal{H}, [\cdot, \cdot])$ is a Hilbert space. By [3, Proposition 25],

$$[a_0, c_0] \subset \sigma_{++}(A) \cup \rho(A).$$

A similar reasoning holds for $[a_0, c_0] \setminus \{b\} \subset \sigma_{--}(A) \cup \rho(A)$ and Theorem 2.2 is proved. \square

Theorem 2.3 *Assume that the operator A is definitizable and the positive half-axis belongs to the resolvent set of A . Then*

$$\infty \in \sigma_{\pi_+}(A) \cup \sigma_{\pi_-}(A) \cup \rho_e(A).$$

Proof:

By [3, Lemma 10] we have $\infty \in \sigma_{\pi_+}(A) \cup \sigma_{\pi_-}(A) \cup \rho_e(A)$ if and only if $\infty \in \sigma_{++}(A) \cup \sigma_{--}(A) \cup \rho_e(A)$. This is equivalent to $\infty \in \sigma_{++}(A-1) \cup \sigma_{--}(A-1) \cup \rho_e(A-1)$. The latter holds if and only if

$$0 \in \sigma_{++}((A-1)^{-1}) \cup \sigma_{--}((A-1)^{-1}) \cup \rho_e((A-1)^{-1}),$$

see [1]. Here, $(A-1)^{-1}$ is a bounded definitizable operator. By the definitizability of A , some interval $(-\infty, -m)$ belongs to $\sigma_{++}(A) \cup \rho(A)$ or to $\sigma_{--}(A) \cup \rho(A)$. Then Theorem 2.3 follows from Theorem 2.2. \square

3 Dampings which are small compared with A_0

Throughout the rest of this paper we make the following assumptions.

(A) The stiffness operator $A_0 : \mathcal{D}(A_0) \subset H \rightarrow H$ is a self-adjoint uniformly positive operator on a Hilbert space H . We define for $s > 0$ the space $H_s = \mathcal{D}(A_0^s)$ equipped with the norm $\|\cdot\|_{H_s} := \|A_0^s \cdot\|_H$ and H_{-s} is the

completion of H with respect to the norm $\|z\|_{H_{-s}} = \|A_0^{-s}z\|_H$. Thus A_0 restricts (extends, respectively) to a bounded operator $A_0 : H_s \rightarrow H_{s-1}$. We use the same notation A_0 to denote this restriction/extension. We denote the inner product on H by $\langle \cdot, \cdot \rangle_H$ or $\langle \cdot, \cdot \rangle$, and the duality pairing on $H_{-s} \times H_s$ by $\langle \cdot, \cdot \rangle_{H_{-s} \times H_s}$. Note that for $(z', z) \in H \times H_s$, we have

$$\langle z', z \rangle_{H_{-s} \times H_s} = \langle z', z \rangle_H.$$

Moreover, for this section, we assume that the operator D is comparable with A_0 in the following sense.

(D1) The damping operator $D : H_{\frac{1}{2}} \rightarrow H_{-\frac{1}{2}}$ is a bounded operator such that $A_0^{-1/2}DA_0^{-1/2}$ is a bounded self-adjoint operator in H and satisfies

$$\langle Dz, z \rangle_{H_{-\frac{1}{2}} \times H_{\frac{1}{2}}} \geq 0, \quad z \in H_{\frac{1}{2}}.$$

The equation (1) is equivalent to the following standard first-order equation

$$\dot{x}(t) = \mathcal{A}x(t),$$

where $\mathcal{A} : \mathcal{D}(\mathcal{A}) \subset H_{\frac{1}{2}} \times H \rightarrow H_{\frac{1}{2}} \times H$, is given by

$$\mathcal{A} = \begin{bmatrix} 0 & I \\ -A_0 & -D \end{bmatrix},$$

$$\mathcal{D}(\mathcal{A}) = \left\{ \begin{pmatrix} z \\ w \end{pmatrix} \in H_{\frac{1}{2}} \times H_{\frac{1}{2}} \mid A_0z + Dw \in H \right\}.$$

The operator \mathcal{A} itself is not self-adjoint in the Hilbert space $H_{\frac{1}{2}} \times H$. It is easy to see (e.g. [34]) that \mathcal{A} has a bounded inverse in $H_{\frac{1}{2}} \times H$ given by

$$\mathcal{A}^{-1} = \begin{bmatrix} -A_0^{-1}D & -A_0^{-1} \\ I & 0 \end{bmatrix}, \quad (3)$$

where $A_0^{-1}D$ is considered as an operator acting in $H_{\frac{1}{2}}$. This together with the fact that

$$J\mathcal{A}, \quad \text{where } J = \begin{bmatrix} I & 0 \\ 0 & -I \end{bmatrix}, \quad (4)$$

is a symmetric operator in the Hilbert space $H_{\frac{1}{2}} \times H$, imply the self-adjointness of $J\mathcal{A}$ in $H_{1/2} \times H$. Therefore, (compare also [33, Proof of Lemma 4.5])

$$\mathcal{A}^* = J\mathcal{A}J, \quad \text{with } \mathcal{D}(\mathcal{A}^*) = J\mathcal{D}(\mathcal{A})$$

and

$$\operatorname{Re} \langle \mathcal{A}x, x \rangle \leq 0 \quad \text{for } x \in \mathcal{D}(\mathcal{A}) \quad \text{and} \quad \operatorname{Re} \langle \mathcal{A}^*x, x \rangle \leq 0 \quad \text{for } x \in \mathcal{D}(\mathcal{A}^*).$$

Hence, \mathcal{A} is the generator of a strongly continuous semigroup of contractions on the state space $H_{\frac{1}{2}} \times H$. This fact is well-known.

For $\begin{pmatrix} x_1 \\ y_1 \end{pmatrix}, \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \in H_{\frac{1}{2}} \times H$ we define an indefinite inner product on $H_{\frac{1}{2}} \times H$ by

$$\left[\begin{pmatrix} x_1 \\ y_1 \end{pmatrix}, \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \right] := \left\langle J \begin{pmatrix} x_1 \\ y_1 \end{pmatrix}, \begin{pmatrix} x_2 \\ y_2 \end{pmatrix} \right\rangle = \langle x_1, x_2 \rangle_{H_{\frac{1}{2}}} - \langle y_1, y_2 \rangle. \quad (5)$$

Then $(H_{\frac{1}{2}} \times H, [\cdot, \cdot])$ is a Krein space and \mathcal{A} is a self-adjoint operator with respect to $[\cdot, \cdot]$. Therefore, we obtain the following proposition.

Proposition 3.1 *The operator \mathcal{A} is self-adjoint in the Krein space $(H_{\frac{1}{2}} \times H, [\cdot, \cdot])$, its spectrum is contained in the closed left half plane and lies symmetric with respect to the real line. The operator \mathcal{A} has a bounded inverse and is the generator of a strongly continuous semigroup of contractions on $H_{\frac{1}{2}} \times H$.*

In the following we will consider the case of an operator A_0 with a compact resolvent. Note that in this case the operator \mathcal{A}^{-1} is, in general, not a compact operator in $H_{\frac{1}{2}} \times H$.

Lemma 3.2 *Assume that the operator A_0^{-1} is compact in H . Then*

$$\sigma_{\text{ess}}(\mathcal{A}) = \left\{ \lambda \in \mathbb{C} \setminus \{0\} \mid \frac{1}{\lambda} \in \sigma_{\text{ess}}(-A_0^{-1}D) \right\} \subset (-\infty, 0). \quad (6)$$

Here $A_0^{-1}D$ is considered as an operator acting in $H_{\frac{1}{2}}$. In particular, the following statements are equivalent.

- (i) *The operator \mathcal{A}^{-1} is a compact operator in $H_{\frac{1}{2}} \times H$.*
- (ii) $\sigma_{\text{ess}}(A_0^{-1}D) = \{0\}$.
- (iii) *The operator D is a compact operator acting from $H_{\frac{1}{2}}$ into $H_{-\frac{1}{2}}$.*

Proof:

Relation (6) follows from (3) (see e.g. [21]). Obviously, by (3), (i) implies (ii). The operator $A_0^{-1}D$ is a bounded self-adjoint operator in the Hilbert space $H_{\frac{1}{2}}$. Therefore, (ii) implies the compactness of $A_0^{-1}D$ in $H_{\frac{1}{2}}$ and, by (3), (i) follows. The operator \mathcal{A}^{-1} considered as an operator acting from $H_{-\frac{1}{2}}$ into $H_{\frac{1}{2}}$ is unitary. This shows the equivalence of (ii) and (iii). \square

Recall, that an at most countably infinite set \mathcal{M} of elements of a Hilbert space is said to be a *Riesz basis* if there exists an isomorphic mapping \mathcal{M} onto an orthonormal basis, cf. [31, Lecture VI]. The following theorem was proved in [22].

Theorem 3.3 *Assume that the operator A_0^{-1} is compact in H and that $0 \notin \sigma_{ess}(A_0^{-1}D)$, where $A_0^{-1}D$ is considered as an operator acting in $H_{\frac{1}{2}}$. Then \mathcal{A} is a definitizable operator and generates an analytic semigroup on $H_{1/2} \times H$. If, in addition, the set $\sigma_{ess}(A_0^{-1}D)$ is countably and has at most countable many accumulation points and if for all $\mu \in \sigma_{ess}(-A_0^{-1}D)$ we have $\frac{1}{\mu} \notin \sigma_p(\mathcal{A})$, then there exists a Riesz basis of $H_{\frac{1}{2}} \times H$ consisting of eigenvectors and finitely many associated vectors of \mathcal{A} .*

The next theorem is the main result of this note. It shows that the the operator \mathcal{A} is definitizable if and only if $0 \notin \sigma_{ess}(A_0^{-1}D)$.

Theorem 3.4 *Assume that the operator A_0^{-1} is compact in H and $0 \in \sigma_{ess}(A_0^{-1}D)$, where $A_0^{-1}D$ is considered as an operator acting in $H_{\frac{1}{2}}$. Then \mathcal{A} is not a definitizable operator in $H_{\frac{1}{2}} \times H$.*

Proof:

Assume that \mathcal{A} is a definitizable operator. As $0 \in \sigma_{ess}(A_0^{-1}D)$, there exists a sequence $(x_n) \in H_{\frac{1}{2}}$ which converges weakly to zero in $H_{\frac{1}{2}}$ with $\|x_n\|_{H_{\frac{1}{2}}} = 1$ and $A_0^{-1}Dx_n \rightarrow 0$ as $n \rightarrow \infty$. Therefore, $(A_0^{\frac{1}{2}}x_n)$ converges weakly to zero in H , hence (x_n) converges to zero in H . The sequence $((\begin{smallmatrix} x_n \\ 0 \end{smallmatrix})) \subset H_{\frac{1}{2}} \times H$ satisfy

$$\|(\begin{smallmatrix} x_n \\ 0 \end{smallmatrix})\|_{H_{\frac{1}{2}} \times H} = 1 \quad \text{and} \quad \mathcal{A}^{-1}(\begin{smallmatrix} x_n \\ 0 \end{smallmatrix}) = \begin{pmatrix} -A_0^{-1}Dx_n \\ x_n \end{pmatrix} \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Moreover, $((\begin{smallmatrix} x_n \\ 0 \end{smallmatrix}))$ converges weakly to zero in $H_{\frac{1}{2}} \times H$ with

$$\limsup_{n \rightarrow \infty} [(\begin{smallmatrix} x_n \\ 0 \end{smallmatrix}), (\begin{smallmatrix} x_n \\ 0 \end{smallmatrix})] = 1.$$

By [3, Lemma 14], 0 is not in $\sigma_{\pi_-}(\mathcal{A}^{-1})$.

We choose a sequence $(y_n) \in H$ which converges weakly to zero in H with $\|y_n\| = 1$. Therefore, $(A_0^{-1}y_n)$ converges to zero in H . The sequence $((\begin{smallmatrix} 0 \\ y_n \end{smallmatrix})) \subset H_{\frac{1}{2}} \times H$ satisfy

$$\|(\begin{smallmatrix} 0 \\ y_n \end{smallmatrix})\|_{H_{\frac{1}{2}} \times H} = 1 \quad \text{and} \quad \mathcal{A}^{-1}(\begin{smallmatrix} 0 \\ y_n \end{smallmatrix}) = \begin{pmatrix} -A_0^{-1}y_n \\ 0 \end{pmatrix} \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Moreover, $((\begin{smallmatrix} 0 \\ y_n \end{smallmatrix}))$ converges weakly to zero in $H_{\frac{1}{2}} \times H$ with

$$\liminf_{n \rightarrow \infty} [(\begin{smallmatrix} 0 \\ y_n \end{smallmatrix}), (\begin{smallmatrix} 0 \\ y_n \end{smallmatrix})] = -11.$$

By [3, Theorem 14], 0 is not in $\sigma_{\pi_+}(\mathcal{A}^{-1})$, hence $0 \notin \sigma_{\pi_+}(\mathcal{A}^{-1}) \cup \sigma_{\pi_-}(\mathcal{A}^{-1})$. It is easily seen (see e.g. [1]) that this is the case if and only if $\infty \notin \sigma_{\pi_+}(\mathcal{A}) \cup \sigma_{\pi_-}(\mathcal{A})$. But this is by Theorem 2.3 impossible. Hence \mathcal{A} is not definitizable. \square

Remark 3.5 *If B is a self-adjoint operator in H such that for some $0 \leq \theta < \frac{1}{2}$ we have*

$$k_1 \|A_0^\theta x\| \leq \|Bx\| \leq k_2 \|A_0^\theta x\| \quad \text{for } x \in \mathcal{D}(B) = \mathcal{D}(A_0^\theta),$$

then the semigroup generated by \mathcal{A} is not analytic, cf. [11, Proposition 2.1]. However, it is shown in [11] that if B is a self-adjoint operator in H with $\mathcal{D}(A_0^{\frac{\theta}{2}}) = \mathcal{D}(B^{\frac{1}{2}})$ for some $\frac{1}{2} \leq \theta \leq 1$ and

$$k_1 A_0^\theta \leq B \leq k_2 A_0^\theta \quad x \in \mathcal{D}(B),$$

then the semigroup generated by \mathcal{A} is analytic. Moreover [25] give sufficient conditions for the Riesz basis property.

References

- [1] T.Ya. Azizov, J. Behrndt, P. Jonas and C. Trunk, “Spectral points of type π_+ and π_- for closed linear relations in Krein spaces”, (submitted).
- [2] T.Ya. Azizov and I.S. Iokhvidov, *Linear Operators in Spaces with an Indefinite Metric*, John Wiley & Sons, Ltd., Chichester, 1989.
- [3] T.Ya. Azizov, P. Jonas and C. Trunk, “Spectral points of type π_+ and π_- of self-adjoint operators in Krein spaces”, *J. Funct. Anal.*, Vol. 226, pg. 114-137, 2005.
- [4] H.T. Banks and K. Ito, “A unified framework for approximation in inverse problems for distributed parameter systems”, *Control Theory and Adv. Tech.*, Vol. 4, pg. 73-90, 1988.
- [5] H.T. Banks, K. Ito and Y. Wang, “Well posedness for damped second order systems with unbounded input operators”, *Differential Integral Equations*, Vol. 8, pg. 587-606, 1995.
- [6] A. Bátkai and K.-J. Engel, “Exponential decay of 2×2 operator matrix semigroups”, *J. Comp. Anal. Appl.*, Vol. 6, pg. 153-164, 2004.

- [7] C.D. Benchimol, “A note on weak stabilizability of contraction semi-groups”, *SIAM J. Control Optimization*, Vol. 16, No. 3, pg. 373-379, 1978.
- [8] J. Bognár, *Indefinite Inner Product Spaces*, Springer Verlag, New York-Heidelberg, 1974.
- [9] S. Chen, K. Liu and Z. Liu, “Spectrum and stability for elastic systems with global or local Kelvin-Voigt damping”, *SIAM J. Appl. Math.*, Vol. 59, No. 2, pg. 651-668, 1998.
- [10] G. Chen and D. Russell, “A mathematical model for linear elastic systems with structural damping”, *Q. Appl. Math.*, Vol. 39, pg. 433-454, 1982.
- [11] S. Chen and R. Triggiani, “Proof of extensions of two conjectures on structural damping for elastic systems”, *Pacific J. Math.*, Vol. 136, No. 1, pg. 15-55, 1989.
- [12] S. Chen and R. Triggiani, “Characterization of domains of fractional powers of certain operators arising in elastic systems, and applications”, *J. Differ. Equations*, Vol. 88, No. 2, pg. 279-293, 1990.
- [13] R.O. Griniv and A.A. Shkalikov, “Operator models in elasticity theory and hydromechanics and the associated analytic semigroups”, *Mosc. Univ. Math. Bull.*, Vol. 54, No. 5, pg. 1-10, 1999.
- [14] E. Hendrickson and I. Lasiecka, “Numerical approximations and regularizations of Riccati equations arising in hyperbolic dynamics with unbounded control operators”, *Comput. Optim. Appl.*, Vol. 2, No. 4, pg. 343-390, 1993.
- [15] E. Hendrickson and I. Lasiecka, “Finite-dimensional approximations of boundary control problems arising in partially observed hyperbolic systems”, *Dynam. Contin. Discrete Impuls. Systems* Vol. 1, No. 1, pg. 101-142, 1995.
- [16] R.O. Hryniv and A.A. Shkalikov, “Exponential stability of semigroups related to operator models in mechanics”, *Math. Notes*, Vol. 73, No. 5, pg. 618-624, 2003.
- [17] R.O. Hryniv and A.A. Shkalikov, “Exponential decay of solution energy for equations associated with some operator models of mechanics”, *Functional Analysis and Its Applications*, Vol. 38, No. 3, pg. 163-172, 2004.

- [18] F. Huang, “On the mathematical model for linear elastic systems with analytic damping”, *SIAM J. Control Optim.*, Vol. 26, No. 3, pg. 714-724, 1988.
- [19] F. Huang, “Some problems for linear elastic systems with damping”, *Acta Math. Sci.*, Vol. 10, No. 3, pg. 319-326, 1990.
- [20] B. Jacob, K. Morris and C. Trunk, “Minimum-phase infinite-dimensional second-order systems”, to appear in *IEEE Transactions on Automatic Control*.
- [21] B. Jacob and C. Trunk, “Location of the spectrum of operator matrices which are associated to second order equations” *Operators and Matrices*, Vol. 1, No. 1, 45-60, 2007.
- [22] B. Jacob, C. Trunk and M. Winklmeier, “Analyticity and Riesz basis property of semigroups associated to damped vibrations” (submitted).
- [23] P. Jonas, “On locally definite operators in Krein spaces”. In *Spectral analysis and its applications*, Theta, Bucharest Vol. 2, pg. 95-127, 2003.
- [24] P. Lancaster, A. Markus and V. Matsaev, “Definitizable operators and quasihyperbolic operator polynomials” *J. Funct. Anal.*, Vol. 131, No. 1, pg. 1-28, 1995.
- [25] P. Lancaster and A. Shkalikov, “Damped vibrations of beams and related spectral problems”, *Canadian Applied Mathematics Quarterly*, Vol. 2, No. 1, pg. 45-90, 1994.
- [26] H. Langer, “Spectral functions of definitizable operators in Krein spaces”. In *Functional Analysis (Dubrovnik, 1981)*, Springer, Berlin, Vol. 948, pg. 1-46, 1982.
- [27] H. Langer, A. Markus and V. Matsaev, “Locally definite operators in indefinite inner product spaces”, *Math. Ann.*, Vol. 308, No. 3, pg. 405-424, 1997.
- [28] I. Lasiecka, “Stabilization of wave and plate equations with nonlinear dissipation on the boundary”, *J. Differential Equations*, Vol. 79, No. 2, pg. 340-381, 1989.
- [29] I. Lasiecka and R. Triggiani, “Uniform exponential energy decay of wave equations in a bounded region with $L^2(0, \infty; L^2(\Gamma))$ -feedback

control in the Dirichlet boundary condition”, *J. Differential Equations*, Vol. 66, No. 3, pg. 340-390, 1987.

- [30] N. Levan, “The stabilization problem: A Hilbert space operator decomposition approach”, *IEEE Trans. Circuits and Systems*, Vol. 25, No. 9, pg. 721-727, 1978.
- [31] N. Nikol’skiĭ *Treatise on the shift operator* Vol. 273 Springer-Verlag, Berlin, 1986.
- [32] M. Slemrod, “Stabilization of boundary control systems”, *J. Differential Equation*, Vol. 22, No. 2, pg. 402-415, 1976.
- [33] M. Tucsnak and G. Weiss, “How to get a conservative well-posed system out of thin air”, Part I, *ESAIM Control Optim. Calc. Var.*, Vol. 9, pg. 247-274, 2003.
- [34] M. Tucsnak and G. Weiss, “How to get a conservative well-posed system out of thin air”, Part II, *SIAM J. Control Optim.*, Vol. 42, No. 3, pg. 907-935, 2003.
- [35] K. Veselić, “Energy decay of damped systems”, *ZAMM*, Vol. 84, pg. 856-864, 2004.

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